

HDG Methods for Elasticity Problems

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Goals:

- Linear elasticity
- Nonlinear elasticity
- Elastodynamics

Notation: Div, Curl, Grad

Divergence (**Div**) of a vector:

$$\nabla \cdot \mathbf{u} = \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3}$$

Rotation (**Curl**) of a vector:

$$\nabla \times \mathbf{u} = \left(\frac{\partial u_3}{\partial x_2} - \frac{\partial u_2}{\partial x_3}, \frac{\partial u_1}{\partial x_3} - \frac{\partial u_3}{\partial x_1}, \frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2} \right)$$

Gradient (**Grad**) of a vector:

$$\mathbf{H} \equiv \nabla \mathbf{u} = \begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{\partial u_1}{\partial x_2} & \frac{\partial u_1}{\partial x_3} \\ \frac{\partial u_2}{\partial x_1} & \frac{\partial u_2}{\partial x_2} & \frac{\partial u_2}{\partial x_3} \\ \frac{\partial u_3}{\partial x_1} & \frac{\partial u_3}{\partial x_2} & \frac{\partial u_3}{\partial x_3} \end{bmatrix}$$

Notation: Strain and Rotation Tensor

Symmetric gradient (strain) tensor:

$$\boldsymbol{\varepsilon} \equiv \frac{\nabla \mathbf{u} + (\nabla \mathbf{u})^T}{2} = \begin{bmatrix} \frac{\partial u_1}{\partial x_1} & \frac{1}{2} \left(\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) & \frac{1}{2} \left(\frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right) \\ \frac{1}{2} \left(\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) & \frac{\partial u_2}{\partial x_2} & \frac{1}{2} \left(\frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right) \\ \frac{1}{2} \left(\frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right) & \frac{1}{2} \left(\frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right) & \frac{\partial u_3}{\partial x_3} \end{bmatrix}$$

Anti-symmetric gradient (rotation) tensor:

$$\boldsymbol{\omega} \equiv \frac{\nabla \mathbf{u} - (\nabla \mathbf{u})^T}{2} = \begin{bmatrix} 0 & \frac{1}{2} \left(\frac{\partial u_1}{\partial x_2} - \frac{\partial u_2}{\partial x_1} \right) & \frac{1}{2} \left(\frac{\partial u_1}{\partial x_3} - \frac{\partial u_3}{\partial x_1} \right) \\ \frac{1}{2} \left(\frac{\partial u_2}{\partial x_1} - \frac{\partial u_1}{\partial x_2} \right) & 0 & \frac{1}{2} \left(\frac{\partial u_2}{\partial x_3} - \frac{\partial u_3}{\partial x_2} \right) \\ \frac{1}{2} \left(\frac{\partial u_3}{\partial x_1} - \frac{\partial u_1}{\partial x_3} \right) & \frac{1}{2} \left(\frac{\partial u_3}{\partial x_2} - \frac{\partial u_2}{\partial x_3} \right) & 0 \end{bmatrix}$$

Note that

$$\mathbf{H} = \boldsymbol{\varepsilon} + \boldsymbol{\omega}.$$

Notation: Stress tensor

Stress tensor:

$$\boldsymbol{\sigma} = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix}$$

For linear elasticity, the stress tensor is related to the strain tensor by

$$\boldsymbol{\sigma} = \boldsymbol{\mathcal{C}}\boldsymbol{\varepsilon}, \quad \boldsymbol{\varepsilon} = \boldsymbol{\mathcal{S}}\boldsymbol{\sigma}$$

where $\boldsymbol{\mathcal{C}}$ is the elastic stiffness (fourth-order) tensor and $\boldsymbol{\mathcal{S}}$ is the elastic compliance (fourth-order) tensor. Assuming that they are symmetric.

Linear Elasticity: Stress Formulation

The governing equations are

$$-\nabla \cdot \boldsymbol{\sigma} = \mathbf{b}, \quad \text{in } \Omega. \quad (1)$$

Here \mathbf{b} is a body force

Since $\boldsymbol{\varepsilon} = \boldsymbol{\mathcal{S}}\boldsymbol{\sigma} = \frac{\nabla \mathbf{u} + \nabla \mathbf{u}^T}{2}$, we obtain

$$\begin{aligned} 2\boldsymbol{\mathcal{S}}\boldsymbol{\sigma} - (\nabla \mathbf{u} + \nabla \mathbf{u}^T) &= \mathbf{0}, & \text{in } \Omega, \\ -\nabla \cdot \boldsymbol{\sigma} &= \mathbf{b}, & \text{in } \Omega. \end{aligned} \quad (2)$$

This is the **stress formulation**.

Linear Elasticity: Strain Formulation

The governing equations are

$$-\nabla \cdot \boldsymbol{\sigma} = \boldsymbol{b}, \quad \text{in } \Omega. \quad (3)$$

Here \boldsymbol{b} is a body force

Since $\boldsymbol{\varepsilon} = \frac{\nabla \boldsymbol{u} + \nabla \boldsymbol{u}^T}{2}$ and $\boldsymbol{\sigma} = \boldsymbol{C}\boldsymbol{\varepsilon}$, we obtain

$$\begin{aligned} 2\boldsymbol{\varepsilon} - (\nabla \boldsymbol{u} + \nabla \boldsymbol{u}^T) &= \mathbf{0}, & \text{in } \Omega, \\ -\nabla \cdot \boldsymbol{C}\boldsymbol{\varepsilon} &= \boldsymbol{b}, & \text{in } \Omega. \end{aligned} \quad (4)$$

This is the **strain formulation**.

Linear Elasticity: Boundary Conditions

Displacement boundary condition:

$$\mathbf{u} = \mathbf{g}_D, \quad \text{on } \Gamma_D. \quad (5)$$

Stress boundary condition:

$$\boldsymbol{\sigma} \cdot \mathbf{n} = \mathbf{g}_N, \quad \text{on } \Gamma_N. \quad (6)$$

The complete governing equations:

$$\begin{aligned} 2\boldsymbol{\mathcal{S}}\boldsymbol{\sigma} - (\nabla\mathbf{u} + \nabla\mathbf{u}^T) &= \mathbf{0}, & \text{in } \Omega, \\ -\nabla \cdot \boldsymbol{\sigma} &= \mathbf{b}, & \text{in } \Omega, \\ \mathbf{u} &= \mathbf{g}_D, & \text{on } \Gamma_D, \\ \boldsymbol{\sigma} \cdot \mathbf{n} &= \mathbf{g}_N, & \text{on } \Gamma_N. \end{aligned} \quad (7)$$

Approximation Spaces

We introduce the following spaces

$$W_h^k = \{w \in L^2(\mathcal{T}_h) : w|_K \in \mathcal{P}_k(K), \forall K \in \mathcal{T}_h\},$$

$$\mathbf{V}_h^k = \{\mathbf{v} \in [L^2(\mathcal{T}_h)]^d : \mathbf{v}|_K \in [\mathcal{P}_k(K)]^d, \forall K \in \mathcal{T}_h\},$$

$$\mathbf{Q}_h^k = \{\mathbf{E} \in [L^2(\mathcal{T}_h)]^{d \times d} : \mathbf{E}|_K \in [\mathcal{P}_k(K)]^{d \times d}, \forall K \in \mathcal{T}_h\},$$

$$\mathbf{S}_h^k = \{\boldsymbol{\zeta} \in [L^2(\mathcal{T}_h)]^{d \times d} : \boldsymbol{\zeta}|_K \in [\mathcal{P}_k(K)]^{d \times d}, \boldsymbol{\zeta} = \boldsymbol{\zeta}^T, \forall K \in \mathcal{T}_h\},$$

$$\mathbf{M}_h^k = \{\boldsymbol{\mu} \in [L^2(\mathcal{E}_h)]^d : \boldsymbol{\mu}|_F \in [\mathcal{P}_k(F)]^d, \forall F \in \mathcal{E}_h\}.$$

Note that \mathbf{S}_h^k is the space of polynomials for symmetric tensors.

HDG Method for Linear Elasticity

We seek $(\boldsymbol{\sigma}_h, \mathbf{u}_h, \hat{\mathbf{u}}_h) \in \mathcal{S}_h^k \times \mathcal{V}_h^k \times \mathcal{M}_h^k$ such that

$$(\mathcal{S}\boldsymbol{\sigma}_h, \boldsymbol{\zeta})_{\mathcal{T}_h} + (\mathbf{u}_h, \nabla \cdot \boldsymbol{\zeta})_{\mathcal{T}_h} - \langle \hat{\mathbf{u}}_h, \boldsymbol{\zeta} \cdot \mathbf{n} \rangle_{\partial \mathcal{T}_h} = 0,$$

$$(\boldsymbol{\sigma}_h, \nabla \mathbf{w})_{\mathcal{T}_h} - \langle \hat{\boldsymbol{\sigma}}_h \mathbf{n}, \mathbf{w} \rangle_{\partial \mathcal{T}_h} = (\mathbf{b}, \mathbf{w})_{\mathcal{T}_h},$$

$$\langle \hat{\boldsymbol{\sigma}}_h \mathbf{n}, \boldsymbol{\mu} \rangle_{\partial \mathcal{T}_h \setminus \partial \Omega} + \langle \hat{\mathbf{u}}_h - \mathbf{g}_D, \boldsymbol{\mu} \rangle_{\Gamma_D} + \langle \hat{\boldsymbol{\sigma}}_h \mathbf{n} - \mathbf{g}_N, \boldsymbol{\mu} \rangle_{\Gamma_N} = 0,$$

(8)

for all $(\boldsymbol{\zeta}, \mathbf{w}, \boldsymbol{\mu}) \in \mathcal{S}_h^k \times \mathcal{V}_h^k \times \mathcal{M}_h^k$, where

$$\hat{\boldsymbol{\sigma}}_h \mathbf{n} = \boldsymbol{\sigma}_h \cdot \mathbf{n} - \boldsymbol{\tau}(\mathbf{u}_h - \hat{\mathbf{u}}_h).$$

(9)

Here $\boldsymbol{\tau}$ is the stabilization tensor.

Nonlinear Elasticity

We consider the static equilibrium for a nonlinear elastic body

$$\mathbf{F} - \nabla \varphi = 0 \quad \text{in } \Omega, \quad (10a)$$

$$-\nabla \cdot \mathbf{P}(\mathbf{F}) = \mathbf{b} \quad \text{in } \Omega \quad (10b)$$

$$\varphi = \mathbf{g}_D, \quad \text{on } \Gamma_D \quad (10c)$$

$$\mathbf{P} \cdot \mathbf{n} = \mathbf{g}_N, \quad \text{on } \Gamma_N, \quad (10d)$$

Here φ is the position vector, \mathbf{F} the deformation gradient, and \mathbf{P} the first Piola-Kirchhoff tensor.

Note that

$$\varphi = \mathbf{u} + \mathbf{x}, \quad \mathbf{F} = \nabla \mathbf{u} + \mathbf{I}. \quad (11)$$

The vector φ describes the deformation of the elastic body Ω .

Hyperelastic materials

For hyperelastic materials, the first Piola-Kirchhoff tensor is defined by

$$\mathbf{P} = \frac{\partial \psi}{\partial \mathbf{F}}, \quad (12)$$

where ψ is a strain energy density function.

Neo-Hookean model:

$$\psi = \frac{\mu}{2} (\text{trace}(\mathbf{F}^T \mathbf{F}) - 3 - 2 \ln(\det(\mathbf{F}))) + \frac{\lambda}{2} (\ln(\det(\mathbf{F})))^2 \quad (13)$$

where μ and λ are the Lamé parameters

$$\mu = \frac{E}{2(1 + \nu)}, \quad \lambda = \frac{E\nu}{(1 + \nu)(1 - 2\nu)}. \quad (14)$$

Here E is Young's modulus and ν is Poisson's ratio.

HDG Method for Nonlinear Elasticity

We seek $(\mathbf{F}_h, \boldsymbol{\varphi}_h, \widehat{\boldsymbol{\varphi}}_h) \in \mathbf{Q}_h^k \times \mathbf{V}_h^k \times \mathbf{M}_h^k$ such that

$$\begin{aligned}(\mathbf{F}_h, \mathbf{E})_{\mathcal{T}_h} + (\boldsymbol{\varphi}_h, \nabla \cdot \mathbf{E})_{\mathcal{T}_h} - \langle \widehat{\boldsymbol{\varphi}}_h, \mathbf{E} \cdot \mathbf{n} \rangle_{\partial \mathcal{T}_h} &= 0, \\(P(\mathbf{F}_h), \nabla \mathbf{w})_{\mathcal{T}_h} - \langle \widehat{P}_h \mathbf{n}, \mathbf{w} \rangle_{\partial \mathcal{T}_h} &= (\mathbf{b}, \mathbf{w})_{\mathcal{T}_h}, \\ \langle \widehat{P}_h \mathbf{n}, \boldsymbol{\mu} \rangle_{\partial \mathcal{T}_h \setminus \partial \Omega} + \langle \widehat{\boldsymbol{\varphi}}_h - \mathbf{g}_D, \boldsymbol{\mu} \rangle_{\Gamma_D} + \langle \widehat{P}_h \mathbf{n} - \mathbf{g}_N, \boldsymbol{\mu} \rangle_{\Gamma_N} &= 0,\end{aligned}\tag{15}$$

for all $(\mathbf{E}, \mathbf{w}, \boldsymbol{\mu}) \in \mathbf{Q}_h^k \times \mathbf{V}_h^k \times \mathbf{M}_h^k$, where

$$\widehat{P}_h \mathbf{n} = P(\mathbf{F}_h) \cdot \mathbf{n} - \boldsymbol{\tau}(\boldsymbol{\varphi}_h - \widehat{\boldsymbol{\varphi}}_h).\tag{16}$$

Here $\boldsymbol{\tau}$ is the stabilization tensor.

The governing equations:

$$\begin{aligned}2\mathcal{S}\boldsymbol{\sigma} - (\nabla \mathbf{u} + \nabla \mathbf{u}^T) &= 0, & \text{in } \Omega, \\ \ddot{\mathbf{u}} - \nabla \cdot \boldsymbol{\sigma} &= \mathbf{b}, & \text{in } \Omega, \\ \mathbf{u} &= \mathbf{g}_D, & \text{on } \Gamma_D, \\ \boldsymbol{\sigma} \cdot \mathbf{n} &= \mathbf{g}_N, & \text{on } \Gamma_N.\end{aligned}\tag{17}$$

Introducing $\mathbf{v} = \dot{\mathbf{u}}$, we obtain

$$\begin{aligned}2\mathcal{S}\dot{\boldsymbol{\sigma}} - (\nabla \mathbf{v} + \nabla \mathbf{v}^T) &= 0, & \text{in } \Omega, \\ \dot{\mathbf{v}} - \nabla \cdot \boldsymbol{\sigma} &= \mathbf{b}, & \text{in } \Omega, \\ \mathbf{v} &= \dot{\mathbf{g}}_D, & \text{on } \Gamma_D, \\ \boldsymbol{\sigma} \cdot \mathbf{n} &= \mathbf{g}_N, & \text{on } \Gamma_N.\end{aligned}\tag{18}$$

HDG Method for Linear Elastodynamics

We seek $(\boldsymbol{\sigma}_h, \mathbf{v}_h, \hat{\mathbf{v}}_h) \in \mathcal{S}_h^k \times \mathcal{V}_h^k \times \mathcal{M}_h^k$ such that

$$\begin{aligned}(\mathcal{S}\dot{\boldsymbol{\sigma}}_h, \boldsymbol{\zeta})_{\mathcal{T}_h} + (\mathbf{v}_h, \nabla \cdot \boldsymbol{\zeta})_{\mathcal{T}_h} - \langle \hat{\mathbf{v}}_h, \boldsymbol{\zeta} \cdot \mathbf{n} \rangle_{\partial \mathcal{T}_h} &= 0, \\(\dot{\mathbf{v}}_h, \mathbf{w})_{\mathcal{T}_h} + (\boldsymbol{\sigma}_h, \nabla \mathbf{w})_{\mathcal{T}_h} - \langle \hat{\boldsymbol{\sigma}}_h \mathbf{n}, \mathbf{w} \rangle_{\partial \mathcal{T}_h} &= (\mathbf{b}, \mathbf{w})_{\mathcal{T}_h}, \\ \langle \hat{\boldsymbol{\sigma}}_h \mathbf{n}, \boldsymbol{\mu} \rangle_{\partial \mathcal{T}_h \setminus \partial \Omega} + \langle \hat{\mathbf{v}}_h - \dot{\mathbf{g}}_D, \boldsymbol{\mu} \rangle_{\Gamma_D} + \langle \hat{\boldsymbol{\sigma}}_h \mathbf{n} - \mathbf{g}_N, \boldsymbol{\mu} \rangle_{\Gamma_N} &= 0,\end{aligned}\tag{19}$$

for all $(\boldsymbol{\zeta}, \mathbf{w}, \boldsymbol{\mu}) \in \mathcal{S}_h^k \times \mathcal{V}_h^k \times \mathcal{M}_h^k$, where

$$\hat{\boldsymbol{\sigma}}_h \mathbf{n} = \boldsymbol{\sigma}_h \cdot \mathbf{n} - \boldsymbol{\tau}(\mathbf{v}_h - \hat{\mathbf{v}}_h).\tag{20}$$

Here $\boldsymbol{\tau}$ is the stabilization tensor.

Nonlinear Elastodynamics

We consider the dynamics for a nonlinear elastic body

$$\mathbf{F} - \nabla \varphi = 0 \quad \text{in } \Omega, \quad (21a)$$

$$\ddot{\varphi} - \nabla \cdot \mathbf{P}(\mathbf{F}) = \mathbf{b} \quad \text{in } \Omega \quad (21b)$$

$$\varphi = \mathbf{g}_D, \quad \text{on } \Gamma_D \quad (21c)$$

$$\mathbf{P} \cdot \mathbf{n} = \mathbf{g}_N, \quad \text{on } \Gamma_N, \quad (21d)$$

Introducing $\mathbf{v} = \dot{\varphi}$, we obtain

$$\begin{aligned} \dot{\mathbf{F}} - \nabla \mathbf{v} &= 0 && \text{in } \Omega, \\ \dot{\mathbf{v}} - \nabla \cdot \mathbf{P}(\mathbf{F}) &= \mathbf{b} && \text{in } \Omega \\ \mathbf{v} &= \dot{\mathbf{g}}_D, && \text{on } \Gamma_D \\ \mathbf{P} \cdot \mathbf{n} &= \mathbf{g}_N, && \text{on } \Gamma_N, \end{aligned} \quad (22)$$

HDG Method for Nonlinear Elastodynamics

We seek $(\mathbf{F}_h, \mathbf{v}_h, \hat{\mathbf{v}}_h) \in \mathbf{Q}_h^k \times \mathbf{V}_h^k \times \mathbf{M}_h^k$ such that

$$(\dot{\mathbf{F}}_h, \mathbf{E})_{\mathcal{T}_h} + (\mathbf{v}_h, \nabla \cdot \mathbf{E})_{\mathcal{T}_h} - \langle \hat{\mathbf{v}}_h, \mathbf{E} \cdot \mathbf{n} \rangle_{\partial \mathcal{T}_h} = 0,$$

$$(\dot{\mathbf{v}}_h, \mathbf{w})_{\mathcal{T}_h} + (\mathbf{P}(\mathbf{F}_h), \nabla \mathbf{w})_{\mathcal{T}_h} - \left\langle \hat{\mathbf{P}}_h \mathbf{n}, \mathbf{w} \right\rangle_{\partial \mathcal{T}_h} = (\mathbf{b}, \mathbf{w})_{\mathcal{T}_h},$$

$$\left\langle \hat{\mathbf{P}}_h \mathbf{n}, \boldsymbol{\mu} \right\rangle_{\partial \mathcal{T}_h \setminus \partial \Omega} + \langle \hat{\mathbf{v}}_h - \dot{\mathbf{g}}_D, \boldsymbol{\mu} \rangle_{\Gamma_D} + \left\langle \hat{\mathbf{P}}_h \mathbf{n} - \mathbf{g}_N, \boldsymbol{\mu} \right\rangle_{\Gamma_N} = 0, \quad (23)$$

for all $(\mathbf{E}, \mathbf{w}, \boldsymbol{\mu}) \in \mathbf{Q}_h^k \times \mathbf{V}_h^k \times \mathbf{M}_h^k$, where

$$\hat{\mathbf{P}}_h \mathbf{n} = \mathbf{P}(\mathbf{F}_h) \cdot \mathbf{n} - \boldsymbol{\tau}(\mathbf{v}_h - \hat{\mathbf{v}}_h). \quad (24)$$

Here $\boldsymbol{\tau}$ is the stabilization tensor.

[NPC11] *High-order implicit hybridizable discontinuous Galerkin methods for acoustics and elastodynamics*, J. Comp. Phys., 230 (2011), pp. 3695–3718.

[NP12] *Hybridizable discontinuous Galerkin methods for partial differential equations in continuum mechanics*, J. of Comp. Phys., 231:5955–5988, 2012.